

What is fluorescence?

FLUORESCENCE is the light emitted by an atom or molecule after a finite duration subsequent to the absorption of electromagnetic energy.

Specifically, the emitted light arises from the transition of the excited species from its first excited electronic singlet level to its ground electronic level. (**usually**)

The development of highly sophisticated fluorescent probe chemistries, new laser and microscopy approaches and site-directed mutagenesis has led to many novel applications of fluorescence in the chemical, physical and life sciences. Fluorescence methodologies are now widely used in the biochemical and biophysical areas, in clinical chemistry and diagnostics and in cell biology and molecular biology.

Why fluorescence?

- its pretty!
- it provides information on the molecular environment
- it provides information on dynamic processes on the nanosecond timescale

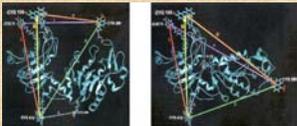
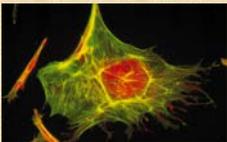
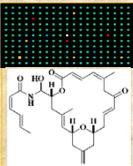
Fluorescence Probes are essentially molecular stopwatches which monitor dynamic events which occur during the excited state lifetime – such as movements of proteins or protein domains

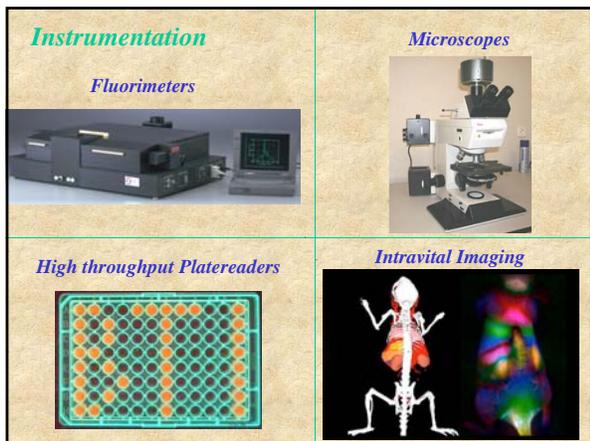
Coupled with modern fluorescence microscopy (confocal, multiphoton, etc) and fluorescent proteins (such as GFP, etc) fluorescence is also providing extremely detailed spatial information in living cells – as well as information on the dynamics of cellular components

Also fluorescence is very, very, very sensitive!

Work with subnanomolar concentrations is routine while femtomolar and even SINGLE MOLECULE studies are possible with some effort

Experimental Systems Accessible to Fluorescence

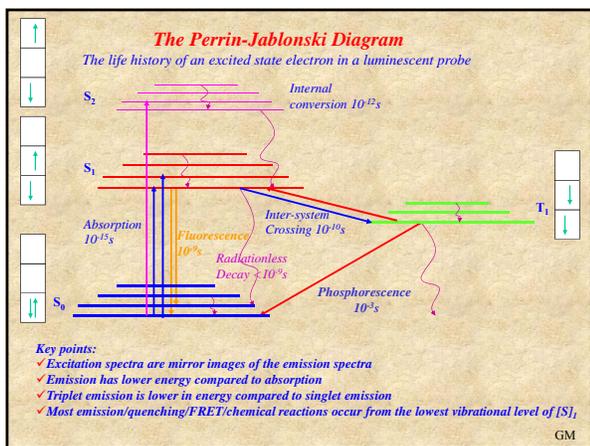
<i>Molecular structure and dynamics</i> 	<i>Cell organization and function</i> 
<i>Live Animals</i> 	<i>Engineered surfaces</i>  <i>High throughput Drug discovery</i>

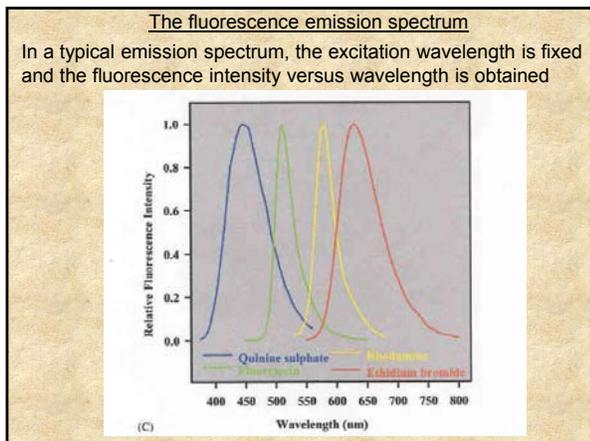


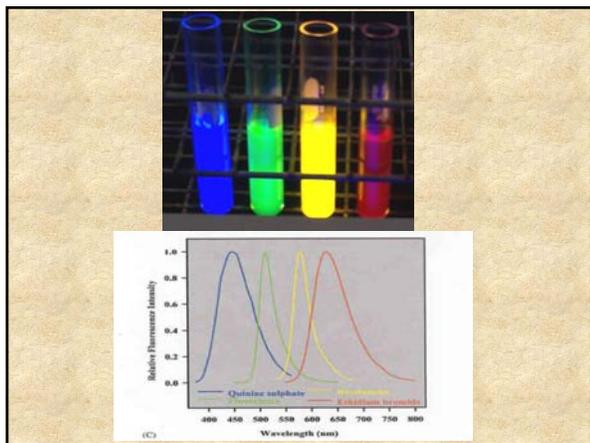
Virtually all fluorescence data required for any research project will fall into one of the following categories.

1. The fluorescence emission spectrum
2. The excitation spectrum of the fluorescence
3. The quantum yield
4. The polarization (anisotropy) of the emission
5. The fluorescence lifetime

In these lectures, we examine each of these categories and briefly discuss historical developments, underlying concepts and practical considerations





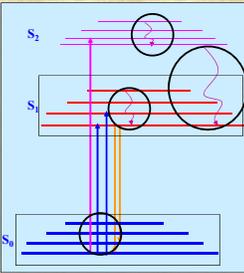


Early examination of a large number of emission spectra resulted in the formulation of certain general rules:

- 1) *In a pure substance existing in solution in a unique form, the fluorescence spectrum is invariant, remaining the same independent of the excitation wavelength*
- 2) *The fluorescence spectrum lies at longer wavelengths than the absorption*
- 3) *The fluorescence spectrum is, to a good approximation, a mirror image of the absorption band of least frequency*

These general observations follow from consideration of the Perrin-Jablonski diagram shown earlier

Specifically, although the fluorophore may be excited into different singlet state energy levels (e.g., S_1 , S_2 , etc) rapid thermalization invariably occurs and emission takes place from the lowest vibrational level of the first excited electronic state (S_1). This fact accounts for the independence of the emission spectrum from the excitation wavelength.



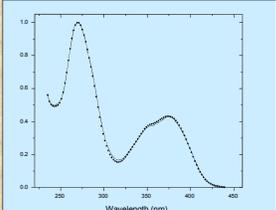
The fact that ground state fluorophores, at room temperature, are predominantly in the lowest vibrational level of the ground electronic state (as required from Boltzmann's distribution law) accounts for the Stokes shift.

Finally, the fact that the spacings of the energy levels in the vibrational manifolds of the ground state and first excited electronic states are usually similar accounts for the fact that the emission and absorption spectra (plotted in energy units such as reciprocal wavenumbers) are approximately mirror images

The fluorescence excitation spectrum

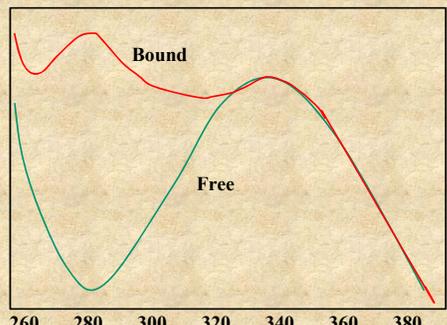
The relative efficiencies of different wavelengths of incident light to excite fluorophores is determined as the excitation spectrum. In this case, the excitation monochromator is varied while the emission wavelength is kept constant if a monochromator is utilized - or the emitted light can be observed through a filter.

If the system is "well-behaved", i.e., if the three general rules outlined above hold, one would expect that the excitation spectrum will match the absorption spectrum. In this case, however, as in the case of the emission spectrum, corrections for instrumentation factors are required.



Overlay of Absorption Spectrum and Corrected Excitation Spectrum for ANS in ethanol

Consider the excitation spectra below for a fluorophore free in solution and then bound to a protein



Quantum Yield

The fluorescence quantum yield was a concept introduced in 1924 by Sergey Ivanovich Vavilov

Perhaps the simplest definition of quantum yield (QY) is that it represents the number of photons emitted divided by the number of photons absorbed:

$$QY = \text{photons emitted as fluorescence} / \text{photons absorbed}$$

So, if a fluorophore absorbs 100 photons and emits 100 photons, its quantum yield is 1.0 or 100% whereas if it absorbs 100 photons and emits only 20 photons then its quantum yield is 0.20 or 20%.

Quantum Yield

A more insightful definition of QY would be that the QY equals the rate of the emission process divided by the sum of the rates of all other deactivation processes.

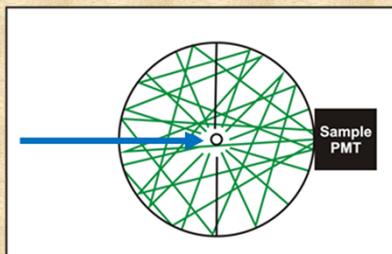
$$QY = k_f / \sum k_d$$

where k_f is the rate of fluorescence, and $\sum k_d$ designates the sum of the rate constants for the various processes, in addition to fluorescence, that depopulate the excited state.

These non-radiative deactivation processes may include photochemical and dissociative processes in which the products are well characterized chemical species (electrons, protons, radicals, molecular isomers). Also, we may have less well characterized changes that result in a return to the ground state with the simultaneous dissipation of the excited state energy as heat. These latter processes are collectively called 'radiationless transitions

The Hamamatsu Corporation recently introduced an instrument for absolute quantum yield determinations, which uses the integrating sphere approach (the C9920-12 External Quantum Efficiency System).

The fluorescing sample is placed in the sphere's interior, which is coated with a highly reflective material such that the emitted light will reflect many times, but will eventually reach the detector. Hence, essentially all of the emitted light can be measured and quantified relative to a highly scattering sample.



The classic approach for determination of a quantum yield using the "relative" method is to select a fluorophore, a standard, which has absorption and emission properties that roughly match those of your sample. For example, tryptophan or N-acetyltryptophanamide (NATA) are common standards for proteins, while quinine sulfate and fluorescein are often used when the excitation is in the mid-UV or visible regions, respectively. Then, one makes up solutions of the standard and sample such that the optical densities are similar. Now, one simply measures the emission spectra of both standard and sample, applies correction factors to get true molecular spectra, and integrates the area under the spectra to get the total intensity of the emitted light. One then uses the following equation to calculate the quantum yield of the sample:

$$QY_{\text{sample}} = QY_{\text{standard}} (I_{\text{T sample}}/I_{\text{T standard}})(1-10^{-OD_{\text{standard}}}/1-10^{-OD_{\text{sample}}})(n_{\text{sample}}/n_{\text{standard}})^2$$

The term $(1-10^{-OD})$ represents the fraction of light absorbed at the excitation wavelength.

I should note that, in the literature, one often sees these terms replaced with absorbances or optical densities, which is only a good approximation if the optical densities are very low. The I_{T} terms are, of course, the total intensities (more about these terms soon), and n refers to the refractive index.

List of quantum yields from "Molecular Fluorescence" by Bernard Valeur

Tab. 6.1. Standards for the determination of fluorescence quantum yields

Range	Compound	Temp. (°C)	Solvent	Φ_{f}	Ref.
270–300 nm	Benzene	20	Cyclohexane	0.05 ± 0.02	1
300–380 nm	Tryptophan	25	H ₂ O (pH 7.2)	0.14 ± 0.02	2
300–400 nm	Naphthalene	20	Cyclohexane	0.23 ± 0.02	3
315–480 nm	2-Aminopyridine	20	0.1 mol L ⁻¹ H ₂ SO ₄	0.60 ± 0.05	4
360–480 nm	Anthracene	20	Ethanol	0.27 ± 0.03	1, 5
400–500 nm	9,10-diphenylanthracene	20	Cyclohexane	0.90 ± 0.02	6, 7
400–600 nm	Quinine sulfate dihydrate	20	0.5 mol L ⁻¹ H ₂ SO ₄	0.546	5, 7
600–650 nm	Rhodamine 101	20	Ethanol	1.0 ± 0.02	8
600–650 nm	Cresyl violet	20	Methanol	0.92 ± 0.02	9
600–650 nm	Cresyl violet	20	Methanol	0.54 ± 0.03	10

- 1) Dawson W. R. and Windsor M. W. (1968) *J. Phys. Chem.* 72, 3251.
- 2) Kirby E. P. and Sietzer R. F. (1970) *J. Phys. Chem.* 74, 4480.
- 3) Berlman I. B. (1965) *Handbook of Fluorescence Spectra of Aromatic Molecules*, Academic Press, London.
- 4) Rosakowicz R. and Tosta A. C. (1968) *J. Phys. Chem.* 72, 2680.
- 5) Melnikow W. H. (1961) *J. Phys. Chem.* 65, 229.
- 6) Hamai S. and Hirayama F. (1983) *J. Phys. Chem.* 87, 83.
- 7) Meech S. R. and Phillips D. (1983) *J. Photochem.* 23, 193.
- 8) Karstens T. and Kubo K. (1980) *J. Phys. Chem.* 84, 1871.
- 9) Audeo-Jacob J., Marx N. J. and Drenthage K. H. (1997) *J. Fluorescence* 7(Suppl), 915.
- 10) Magde D., Brannon J. H., Cramers T. L. and Olmsted J. III (1979) *J. Phys. Chem.* 83, 696.

But you should be aware that quantum yields are notoriously difficult to measure. For example, "reliable" literature values for quinine sulfate range from ~0.50 – 0.70!

Polarization



Polarizers have been in use for a very long time - the Vikings are believed to have used a "sunstone" (thought to have been composed either of the mineral cordierite or iceland spar – calcite – both of which are naturally polarizing materials) to observe the location of the sun on foggy or overcast days. Since scattered sunlight is highly polarized compared to light coming along the direction to the sun, the distribution of the sky's brightness could be observed through the sunstone and hence the sun's position could be localized and, if the time of day were known, the compass directions.

Etienne-Louis Malus (1775-1812)



In 1808, Malus observed sunlight reflected from the windows of the Luxemburg Palace in Paris through an Iceland spar (Calcite) crystal that he rotated.




(Erasmus Bartholin (1625-1698) discovered the double refraction of light by Iceland spar in 1669)

Malus discovered that the intensity of the reflected light varied as he rotated the crystal and coined the term **"polarized"** to describe this property of light.

He published his findings in 1809:
 "Sur une propriété de la lumière réfléchie par les corps diaphanes" (*Bull. Soc. Philomat.* 1:16)

Malus also derived an expression for calculating the transmission of light as a function of the angle (θ) between two polarizers. This equation (Malus' Law) is now written as: $I_a = I_0 (\cos^2\theta)$

Sir David Brewster (1781-1868)



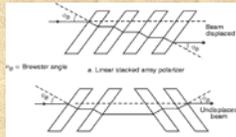
David Brewster studied the relationship between refractive index and angle of incidence on the polarization of the reflected light

III. On the law of the partial polarization of light by reflexion. By DAVID BREWSTER, LL.D. F.R.S. L. & E.
 Read February 4, 1830.

He discovered that for normal glass and visible light, an incidence angle of ~56 degrees resulted in total reflection of one plane of polarization – this angle is now known as *Brewster's Angle*

$$\theta_B = \tan^{-1} \left(\frac{n_2}{n_1} \right)$$

This discovery allowed Brewster to construct a polarizer composed of a "pile of plates"



William Nicol (1770-1851)

In 1828, Nicol joined two crystals of Iceland spar, cut at an of 68° angle, using Canada balsam.

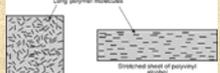


Other important calcite polarizers developed around this time include: Glan-Foucault; Glan-Thompson; Glan-Taylor; Wollaston; Rochon

But the Henry Ford of polarizers was.....

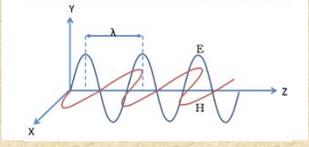
Edwin Herbert Land (1909-1991)

In 1929 Edwin Land patented the sheet polarizer (the J-sheet), consisting of crystals of iodoquinine sulfate embedded in nitrocellulose film followed by alignment of the crystals by stretching which led to dichroism. In 1938 he invented the H-sheet which was comprised of polyvinyl alcohol sheets with embedded iodine.



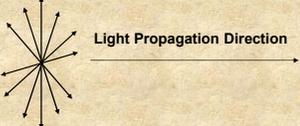
Polarization

Light can be considered as oscillations of an electromagnetic field – characterized by electric and magnetic components - perpendicular to the direction of light propagation.

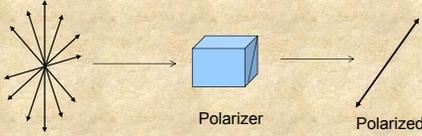


In these lectures we shall be concerned only with the electric component. In natural light the electric field vector can assume any direction of oscillation perpendicular or normal to the light propagation direction.

Unpolarized (natural) light



Polarizers are optically active devices that can isolate one direction of the electric vector.

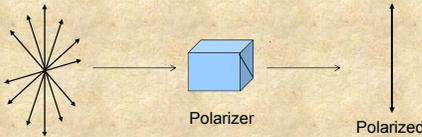


Unpolarized (natural) light

Polarizer

Polarized light

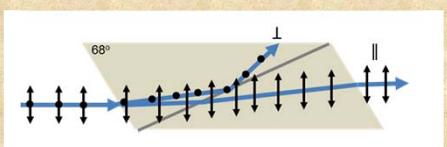
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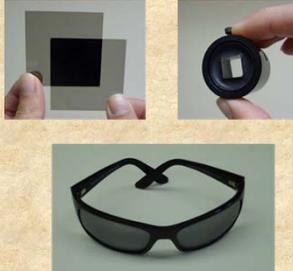
Unpolarized (natural) light

Polarizer

Polarized light



The most common polarizers used today are (1) dichroic devices, which operate by effectively absorbing one plane of polarization (e.g., Polaroid type-H sheets based on stretched polyvinyl alcohol impregnated with iodine) and (2) double refracting calcite (CaCO₃) crystal polarizers - which differentially disperse the two planes of polarization (examples of this class of polarizers are Nicol polarizers, Wollaston prisms and Glan-type polarizers such as the Glan-Foucault, Glan-Thompson and Glan-Taylor polarizers)



Polarization of Fluorescence
 In 1920, F. Weigert discovered that the fluorescence from solutions of dyes was polarized. Specifically, he looked at solutions of fluorescein, eosin, rhodamine and other dyes and noted the effect of temperature and viscosity on the observed polarization.
 In Weigert's words "Der Polarisationsgrad des Fluorezenzlichtes nimmt mit wachsender Molekulargröße, mit zunehmender Viskosität des Mediums und mit abnehmender Temperatur, also mit Verringerung der Beweglichkeit der Einzelteilchen zu"
"The degree of the polarization increases with increasing molecular size, with increasing viscosity of the medium and with decreasing temperature, that is with the reduction of the mobility of the single particles." He recognized that all of these considerations meant that fluorescence polarization increased as the mobility of the emitting species decreased.

In 1925 - 1926, Francis Perrin published several important papers describing a quantitative theory of fluorescence polarization including what is now considered his classic paper containing most of the essential information that we use to this day

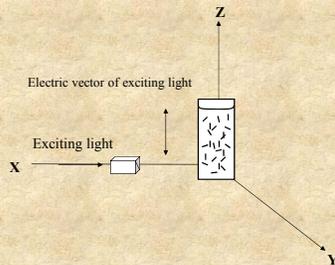
J. de Physique 1926

POLARISATION DE LA LUMIÈRE DE FLUORESCENCE. VIE MOYENNE DES MOLÉCULES
 DANS L'ÉTAT EXCITÉ
 par M. FRANCIS PERRIN.

$$P = P_0 \frac{1}{1 + \left(1 - \frac{1}{3} P_0\right) \frac{RT}{\eta v_f \tau}}$$

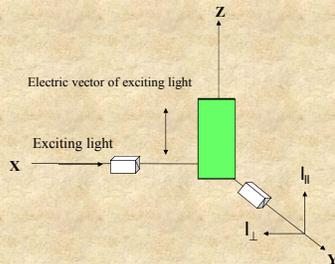
Polarization remained largely in the province of the physicists for almost two decades, until Gregorio Weber began his thesis.
 Weber's subsequent theoretical and experimental work - which extended Perrin's earlier contributions and also developed what became modern instrumentation - brought fluorescence polarization to the attention of the biochemical community, and so ushered in a new scientific discipline - quantitative biological fluorescence.

Consider an XYZ coordinate framework with a fluorescent solution placed at the origin, as shown below, where XZ is in the plane of the page.



In this system, the exciting light is traveling along the X direction. If a polarizer is inserted in the beam, one can isolate a unique direction of the electric vector and obtain light polarized parallel to the Z axis which corresponds to the vertical laboratory axis.

This exciting light will be absorbed by the fluorophore at the origin and give rise to fluorescence which is typically observed at 90° to the excitation direction, i.e., from along the Y axis.



The actual direction of the electric vector of the emission can be determined by viewing the emission through a polarizer which can be oriented alternatively in the parallel or perpendicular direction relative to the Z axis or laboratory vertical direction.

Polarization is then defined as a function of the observed parallel ($I_{||}$) and perpendicular intensities (I_{\perp}):

$$P = \frac{I_{||} - I_{\perp}}{I_{||} + I_{\perp}}$$

If the emission is completely polarized in the parallel direction, i.e., the electric vector of the exciting light is totally maintained, then:

$$P = \frac{1 - 0}{1 + 0} = 1$$

If the emitted light is totally polarized in the perpendicular direction then:

$$P = \frac{0-1}{0+1} = -1$$

The limits of polarization are thus +1 to -1

Another term frequently used in the context of polarized emission is anisotropy (usually designated as either A or r) which is defined as:

$$r = \frac{I_{||} - I_{\perp}}{I_{||} + 2I_{\perp}}$$

By analogy to polarization, the limits of anisotropy are +1 to -0.5.

A comment about the difference between polarization and anisotropy:

Given the definition of polarization and anisotropy, one can show that:

$$r = \frac{2}{3} \left(\frac{1}{P} - \frac{1}{3} \right)^{-1} \quad \text{or} \quad r = \frac{2P}{3-P}$$

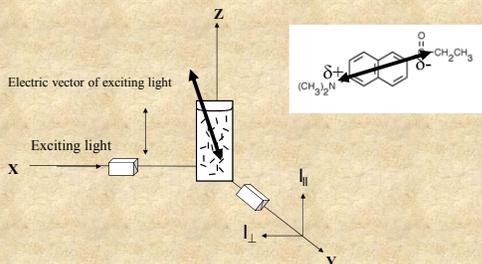
For example:

P	r
0.50	0.40
0.25	0.20
0.10	0.069
-0.10	-0.065

Clearly, the information content in the polarization function and the anisotropy function is identical and the use of one term or the other is dictated by practical considerations as will be discussed later.

In solution these limits (e.g., +/-1) are not realized. Consider, as shown below, fluorophores at the origin of our coordinate system.

Upon absorption of an exciting photon a dipole moment is created in the fluorophore (usually of different magnitude and direction from the ground state dipole). The orientation of this dipole moment relative to the nuclear framework, and its magnitude, will be determined by the nature of the substituents on the molecule. This excited state dipole moment is also known as the transition dipole or transition moment.



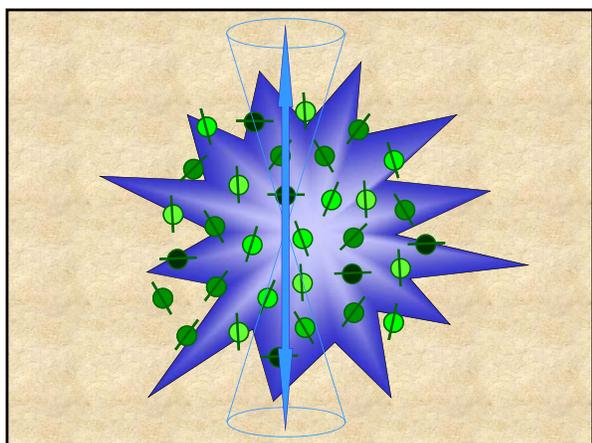
In fact, if light of a particular electric vector orientation (plane polarized light) impinges on a sample, only those molecules which are properly oriented relative to this electric vector can absorb the light.

Specifically, the probability of the absorption is proportional to the cosine squared ($\cos^2\theta$) of the angle θ between the exciting light and the transition dipole.

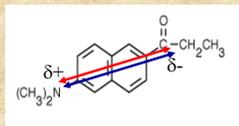
The diagram shows a vertical blue arrow labeled 'Exciting light' and a red arrow labeled 'Absorption Dipole' at an angle θ to the vertical. To the right is a graph titled $\cos^2 \theta$ with a y-axis from 0.00 to 1.00 and an x-axis from 0° to 90° . A red curve starts at (0, 1.00) and ends at (90, 0.00), passing through (45, 0.50).

Hence, when we excite an ensemble of randomly oriented fluorophores with plane-polarized light we are performing a *photoselection* process, creating a population of excited molecules which nominally have their excited dipoles lined up with the polarization direction of the excitation. This process is illustrated below:

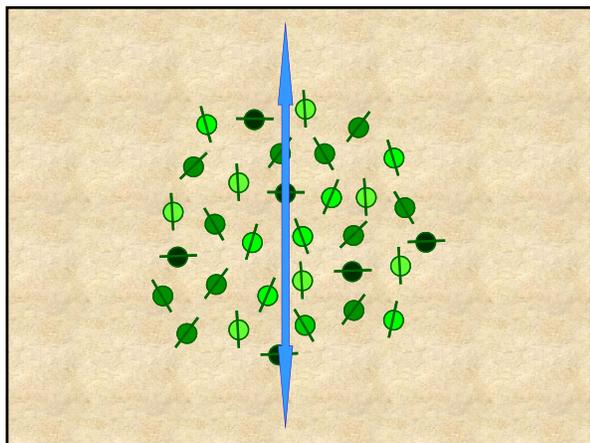
The diagram shows two cylindrical containers. The left one is labeled 'Potential dipoles' and contains randomly oriented black dashes. The right one is labeled 'Excited state dipoles' and contains black dashes all aligned vertically. A wavy arrow labeled $h\nu$ points from the left container to the right.



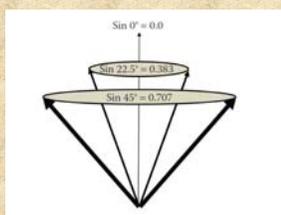
Consider now that the transition dipole corresponding to the emission of light from the excited fluorophore is parallel to the absorption dipole and that the excited fluorophore cannot rotate during the lifetime of the excited state (for example if the fluorophores are embedded in a highly viscous or frozen medium).



If we were to now measure the polarization of the emission it would be less than +1 since some of the dipoles excited will not be exactly parallel to the direction of the exciting light.

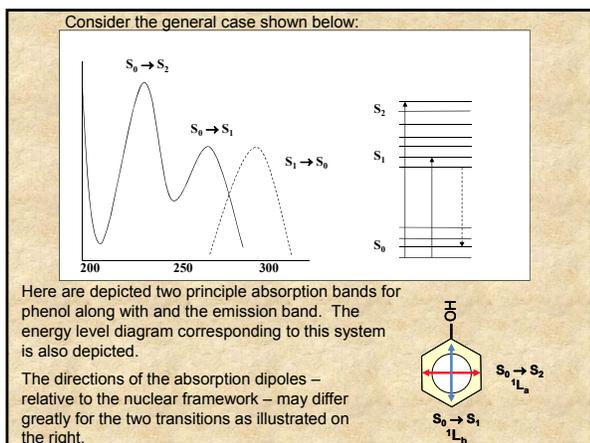


In fact, the number of potential dipoles making an angle θ with the vertical axis will be proportional to $\sin \theta$.



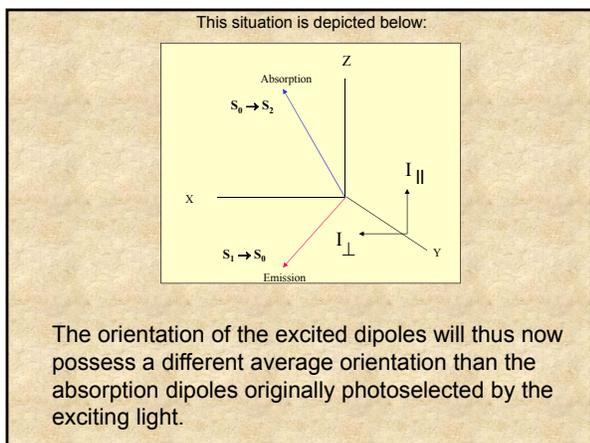
We can then calculate that the upper polarization limit for such a randomly oriented (but rigidly fixed, i.e., non-rotating) ensemble - with co-linear excitation and emission dipole - will be +1/2

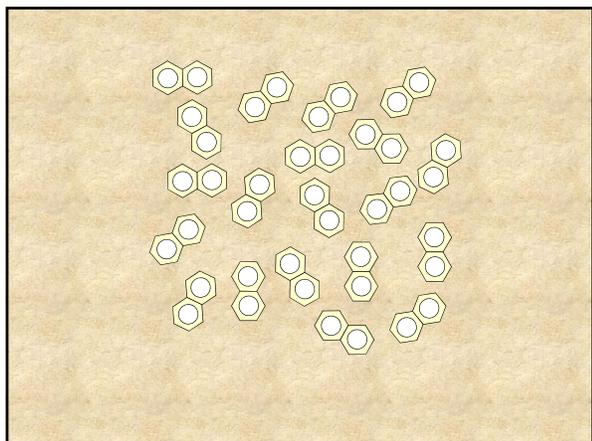
This case, however, assumes that the emission dipole is parallel (co-linear) to the absorption dipole.

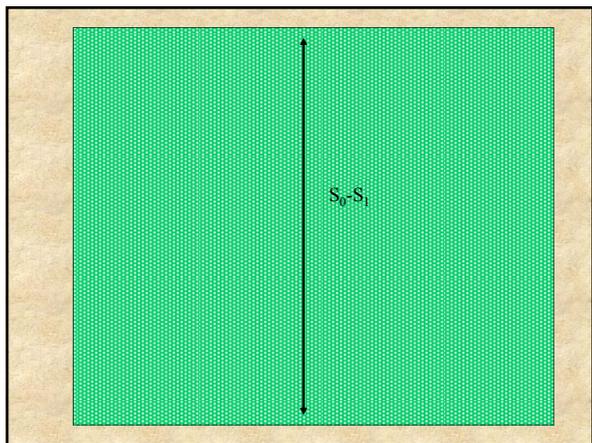


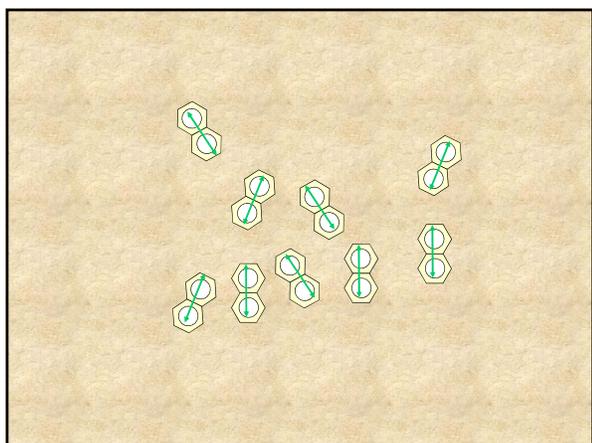
So we see that the two excited dipoles corresponding to the $S_0 \rightarrow S_1$ and the $S_0 \rightarrow S_2$ transitions may be oriented at an arbitrary angle - in the extreme case this angle could be 90° .

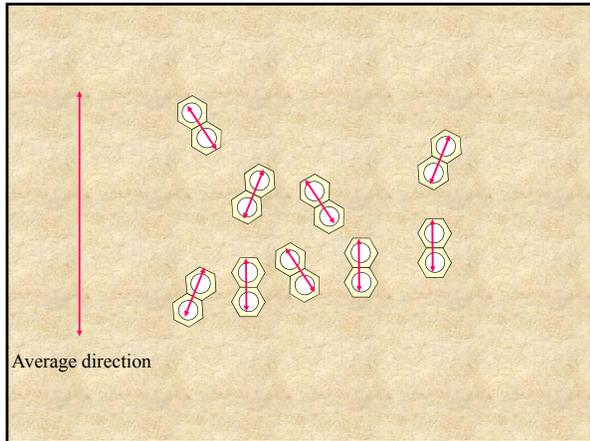
After the excitation process, however, regardless of whether the absorption process corresponded to the $S_0 \rightarrow S_1$ or the $S_0 \rightarrow S_2$ transition, rapid thermalization leaves the excited fluorophore in the S_1 level.

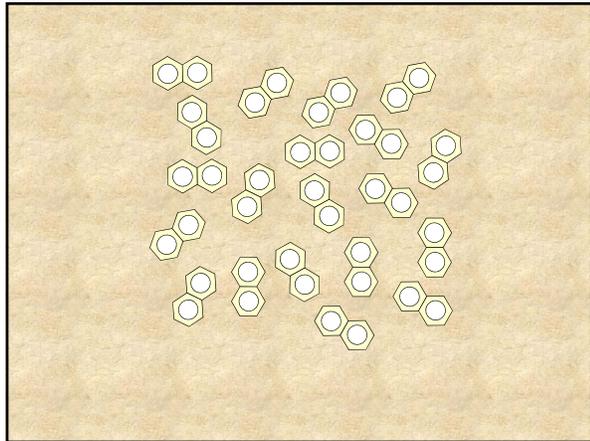


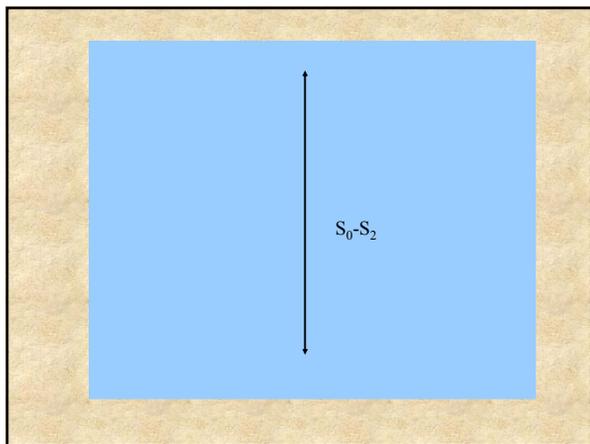


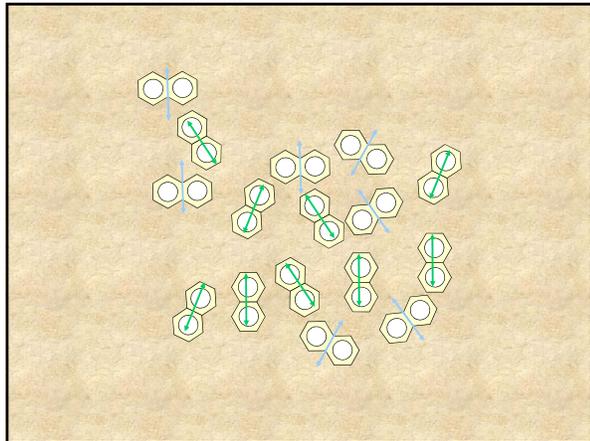


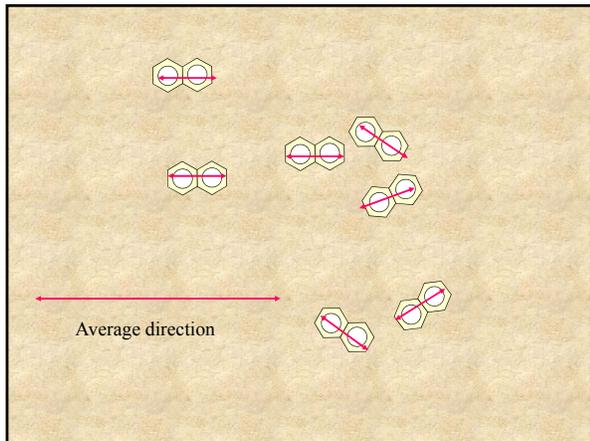












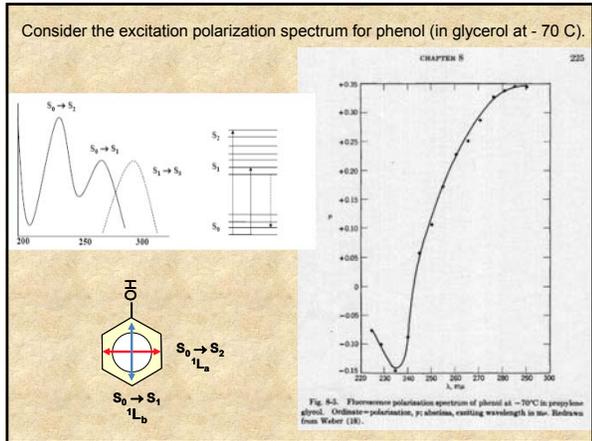
Hence we will observe more emission in the perpendicular direction than in the parallel direction and the resulting polarization will be negative. Considering the same $\cos^2 \theta$ photoselection rule and the $\sin \theta$ population distribution as before we can show that, if the absorption and emission dipoles are at 90° to each other, then $P = -1/3$.

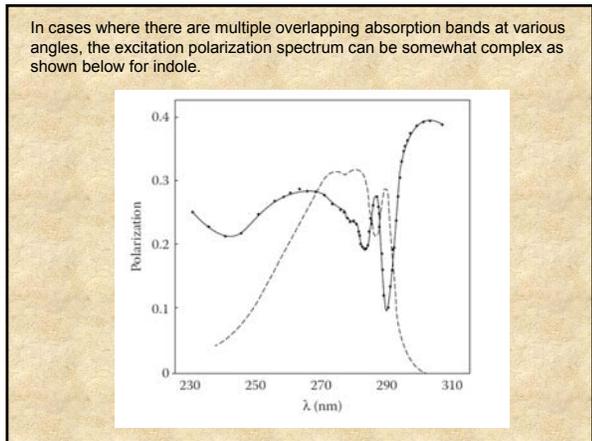
These polarization values, in the absence of rotation, are termed limiting or intrinsic polarizations and are denoted as P_o . In general:

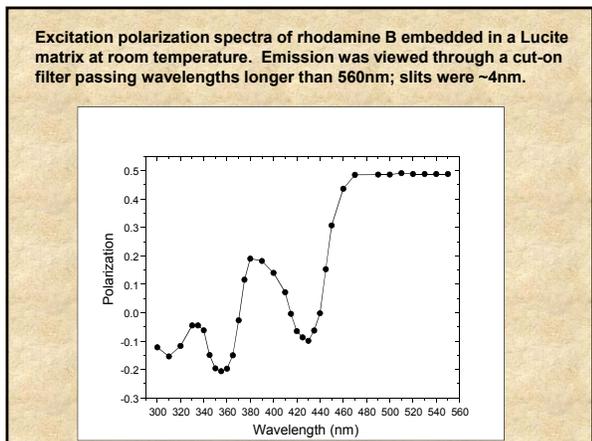
$$\frac{1}{P_o} - \frac{1}{3} = \frac{5}{3} \left(\frac{2}{3 \cos^2 \phi - 1} \right)$$

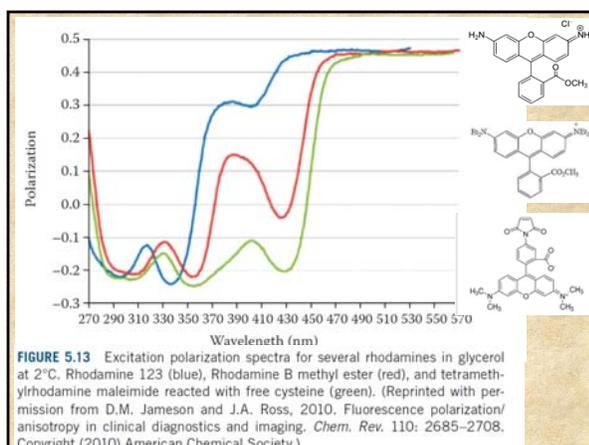
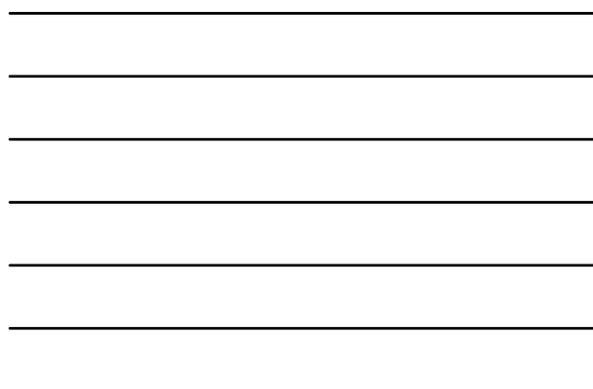
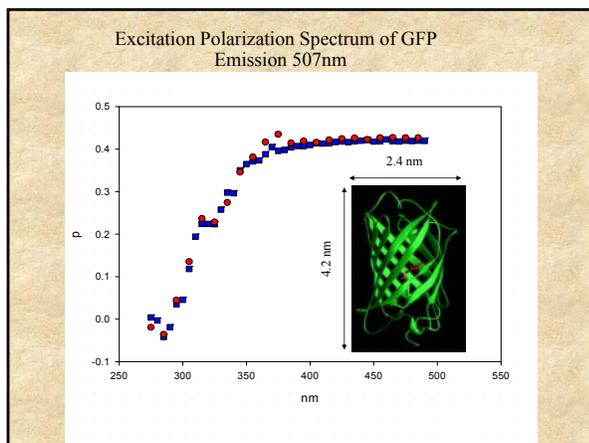
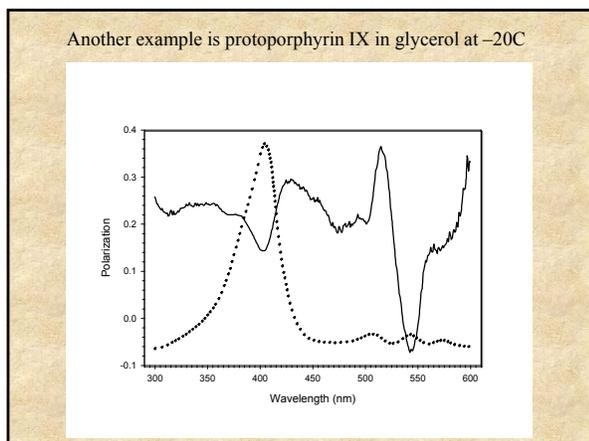
Where ϕ is the angle between absorption and emission dipoles.

We can then understand that the limiting polarization of a fluorophore will depend upon the excitation wavelength.









Note: in the case of multi-photon excitation the limits differ

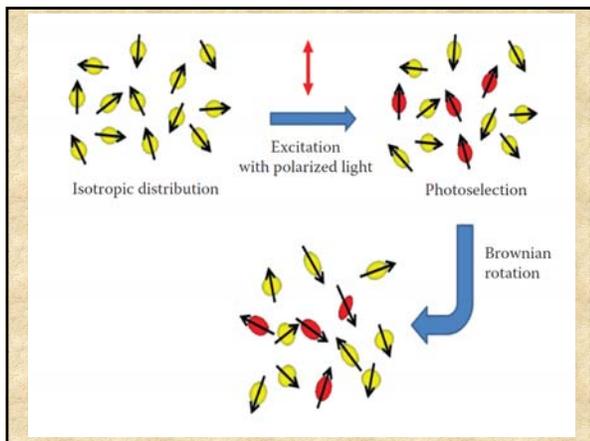
We may now consider the case where the fluorophore is permitted to rotate during the excited state lifetime.

Additional depolarization occurs if the dipole rotates through an angle ω .

In fact:

$$\frac{1}{P} - \frac{1}{3} = \left(\frac{1}{P_0} - \frac{1}{3} \right) \left(\frac{2}{3 \cos^2 \omega - 1} \right)$$

where P is the observed polarization. So the total depolarization is determined by an intrinsic factor (P_0) and an extrinsic factor (ω).



F. Perrin related the observed polarization to the excited state lifetime and the rotational diffusion of a fluorophore: *Perrin, F. 1926. Polarisation de la Lumiere de Fluorescence. Vie Moyene des Molecules Fluorescentes. J. Physique. 7:390-401.*

Specifically:
$$\frac{1}{P} - \frac{1}{3} = \left(\frac{1}{P_0} - \frac{1}{3} \right) \left(1 + \frac{RT}{\eta V} \tau \right)$$

where V is the molar volume of the rotating unit, R is the universal gas constant, T the absolute temperature, η the viscosity and τ the excited state lifetime.

We can rewrite this equation as:
$$\frac{1}{P} - \frac{1}{3} = \left(\frac{1}{P_0} - \frac{1}{3} \right) \left(1 + \frac{3\tau}{\rho} \right)$$

Where ρ is the Debye rotational relaxation time which is the time for a given orientation to rotate through an angle given by the arccos e^{-1} (68.42°).

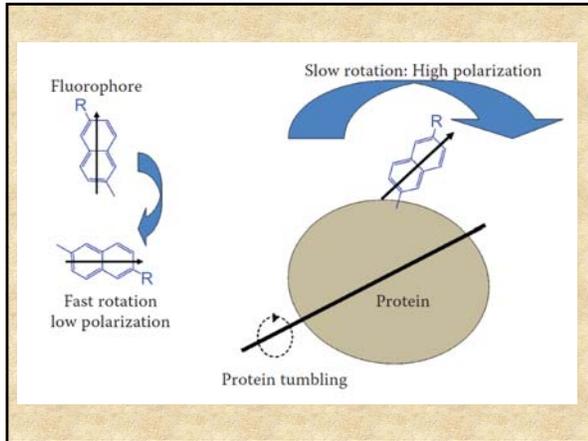
For a spherical molecule:

$$\rho_0 = \frac{3\eta V}{RT}$$

For a spherical protein, it follows that:

$$\rho_0 = \frac{3\eta M(v + h)}{RT}$$

Where M is the molecular weight, v is the partial specific volume and h the degree of hydration.



** Rotational relaxation time versus rotational correlation time.*

We should note that it is not uncommon to see the term "rotational correlation time", often denoted as τ_c , used in place of the Debye rotational relaxation time. The information content of these terms is similar since $\rho = 3\tau_c$ but we have observed that some people become rather fervently attached to the use of one term or the other.

In the original development of the theories of rotational motion of fluorophores Perrin and others used the rotational relaxation time, as originally defined by Debye in his studies on dielectric phenomena. Only later (in the 1950's) during the development of nuclear magnetic resonance was the term rotational correlation time used by Bloch. It thus seems reasonable for fluorescence practitioners to use ρ but certainly adoption of either term should not lead to confusion. In terms of anisotropy and rotational correlation times, then, the Perrin equation would be:

$$\frac{r_0}{r} = \left(1 + \frac{\tau}{\tau_c} \right)$$

In the case of fluorescence probes associated non-covalently with proteins, (for example porphyrins, FAD, NADH or ANS to give but a few systems), the probe is held to the protein matrix by several points of attachment and hence its "local" mobility, that is, its ability to rotate independent of the overall "global" motion of the protein, is very restricted.

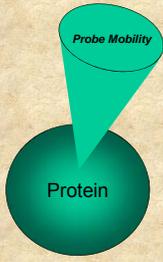
In the case of a probe attached covalently to a protein, via a linkage through an amine or sulfhydryl groups for example, or in the case of tryptophan or tyrosine sidechains, considerable "local" motion of the fluorophore can occur. In addition, the protein may consist of flexible domains which can rotate independent of the overall "global" protein rotation. This type of mobility hierarchy is illustrated on the right for the case of a probe covalently attached to a dimeric protein

Rotational Modalities

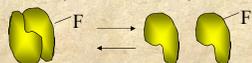
(a) overall dimer rotation
 (b) movement of one C-domain relative to other domains
 (c) movement of dye molecule around its point of attachment

In the case of fluorescence probes associated non-covalently with proteins, (for example porphyrins, FAD, NADH or ANS to give but a few systems), the probe is held to the protein matrix by several points of attachment and hence its "local" mobility, that is, its ability to rotate independent of the overall "global" motion of the protein, is very restricted.

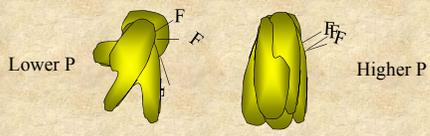
In the case of a probe attached covalently to a protein, via a linkage through an amine or sulfhydryl groups for example, or in the case of tryptophan or tyrosine sidechains, considerable "local" motion of the fluorophore can occur. In addition, the protein may consist of flexible domains which can rotate independent of the overall "global" protein rotation. This type of mobility hierarchy is illustrated on the right for the case of a probe covalently attached to a dimeric protein



Polarization methods are ideally suited to study the aggregation state of a protein. Consider, for example the case of a protein dimer - monomer equilibrium.



Following either intrinsic protein fluorescence (if possible) or by labeling the protein with a suitable probe one would expect the polarization of the system to decrease upon dissociation of the dimer into monomers since the smaller monomers will rotate more rapidly than the dimers (during the excited state lifetime).



Hence for a given probe lifetime the polarization (or anisotropy) of the monomer will be less than that of the dimer

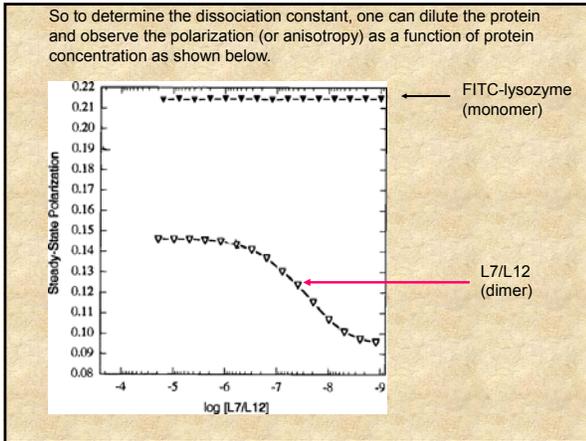
In the concentration range near the dimer/monomer equilibrium constant, one expects to observe a polarization intermediate between that associated with either dimer or monomer. One can relate the observed polarization to the fraction of dimer or monomer using the additivity of polarizations first described by Weber (1952) namely:

$$\left(\frac{1}{\langle P \rangle} - \frac{1}{3} \right)^{-1} = \sum f_i \left(\frac{1}{P_i} - \frac{1}{3} \right)^{-1}$$

where $\langle P \rangle$ is the observed polarization, f_i is the fractional intensity contributed by the i th component and P_i is the polarization of the i th component. One must then relate the fractional intensity contributions to molar quantities which means that one must take into account any change in the quantum yield of the fluorophore associated with either species.

The anisotropy function is directly additive (owing to the fact that the denominator represents the total emitted intensity) and hence:

$$\langle r \rangle = \sum f_i r_i$$



The polarization/anisotropy approach is also very useful to study protein-ligand interactions in general.

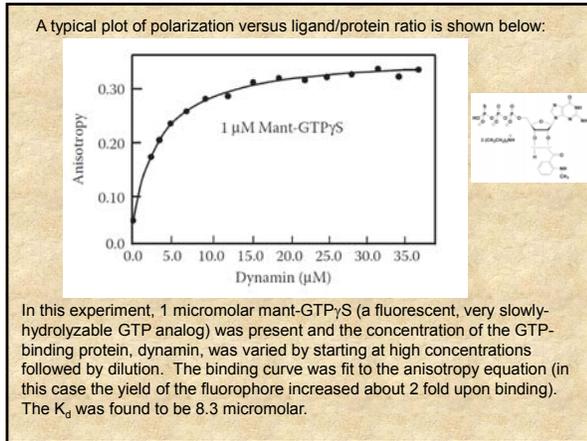
The first application of fluorescence polarization to monitor the binding of small molecules to proteins was carried out by D. Laurence in 1952 using Gregorio Weber's instrumentation in Cambridge. Specifically, Laurence studied the binding of numerous dyes, including fluorescein, eosin, acridine and others, to bovine serum albumin, and used the polarization data to estimate the binding constants.

Although many probes (such as fluorescein) do not significantly alter their quantum yield upon interaction with proteins, one should not take this fact for granted and would be well advised to check. If the quantum yield does in fact change, one can readily correct the fitting equation to take the yield change into account. In terms of anisotropy the correct expression relating observed anisotropy (r) to fraction of bound ligand (x), bound anisotropy (r_b), free anisotropy (r_f), and the quantum yield enhancement factor (g) is:

$$x = \frac{r - r_f}{r_b - r_f + (g - 1)(r_b - r)}$$

In terms of polarization, the analogous equation is

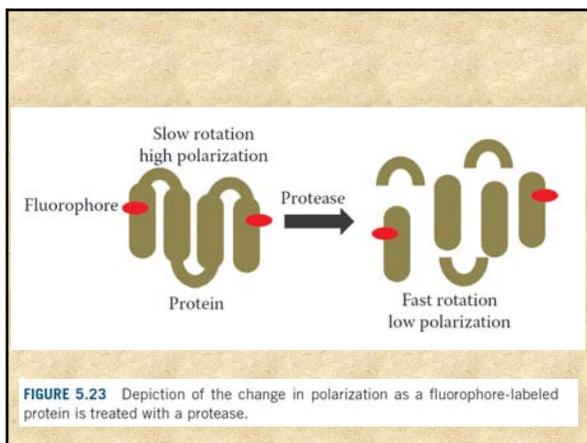
$$x = \frac{(3 - P_b)(P - P_f)}{(3 - P)(P_b - P_f) + (g - 1)(3 - P_f)(P_b - P)}$$



Proteolytic processing, mediated by proteolytic enzymes, or proteases, is critical to many vital biological processes, including post-translational protein processing, blood clotting, digestion, hormone processing, apoptosis, and many others, as well as many deleterious processes, such as those mediated by anthrax and botulinum neurotoxins.

Hence, an evaluation of protease activity is often a requirement for an understanding of a particular pathway or for development of novel therapeutic agents. Protease assays have been around for many decades but more recently the development of rapid and sensitive protease assays suitable for high-throughput screening has attracted considerable attention.

Fluorescence polarization lends itself very well to such assays since the essential aspect of a protease is to cleave a peptide bond which almost always results in smaller molecular weight species. Hence, if the target protein can be labeled with a fluorescence probe one would expect the polarization to decrease after proteolysis since the fluorophore will be able to rotate more rapidly after the protein mass to which it is tethered is reduced in size



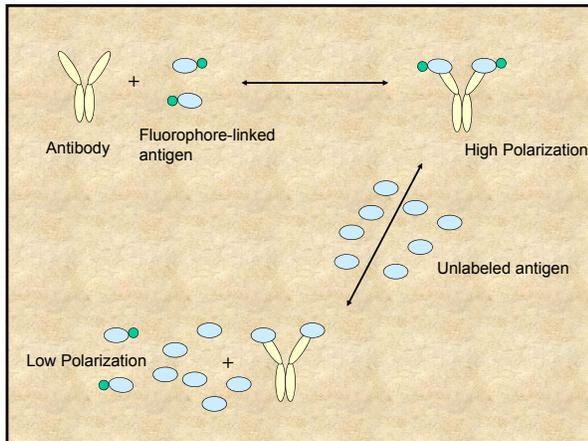
FPIA – Fluorescence Polarization ImmunoAssay

Among the first commercial instruments designed to use a fluorescence polarization immunoassay for clinical diagnostic purposes was the Abbott TDx – introduced in 1981.



The basic principle of a polarization immunoassay is to:

- (1) Add a fluorescent analog of a target molecule – e.g., a drug – to a solution containing antibody to the target molecule
- (2) Measure the fluorescence polarization, which corresponds to the fluorophore bound to the antibody
- (3) Add the appropriate biological fluid, e.g., blood, urine, etc., and measure the decrease in polarization as the target molecules in the sample fluid bind to the antibodies, displacing the fluorescent analogs.



7. Dandliker, W. B., Kelly, R. J., Dandliker, J., et al., Fluorescence polarization immunoassay. Theory and experimental method. *Immunochemistry* 10, 219–227 (1973).

CLIN. CHEM. 27:7, 1190–1197 (1981)

Fluorescence Polarization Immunoassay I. Monitoring Aminoglycoside Antibiotics in Serum and Plasma

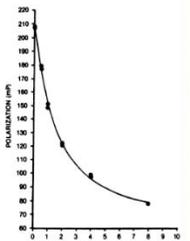
Michael E. Jolley, Stephen D. Stroupe, Chao-Huei J. Wang, Helen N. Panas, Candace L. Keegan, Robert L. Schmidt, and Kathryn S. Schwenzler

Special Apparatus

A microprocessor-controlled fluorometer was constructed in our laboratories (19). The instrument was designed to measure precisely the polarization of fluorescence emitted from the sample contained in a standard 12 × 75 mm disposable culture tube. Each determination was performed in 10 s, to a precision of ±0.001 polarization unit. The instrument determined fluorescence polarization (in arbitrary units) according to the equation (7):

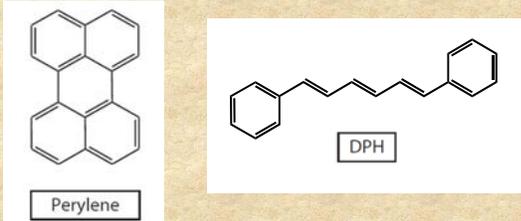
$$P = \frac{I_{\text{parallel}} - I_{\text{perpendicular}}}{I_{\text{parallel}} + I_{\text{perpendicular}}}$$

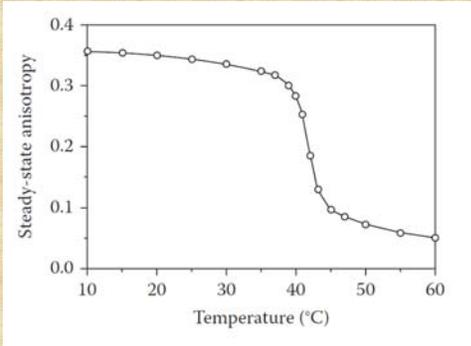
where I represents light intensity. For convenience, results are reported in millipolarization units (mP), where 1 mP = 0.001 P.



Concentration (mg/L)	Polarization (mP)
0	220
1	190
2	170
3	155
4	145
5	135
6	125
7	115
8	105
9	95
10	85

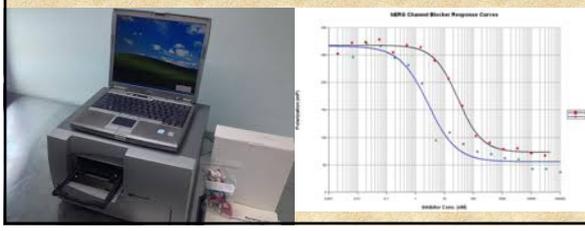
In 1971, Gregorio Weber and his colleagues introduced the use of polarization to study the physical state of lipids in model membrane systems (an idea he had actually mentioned in his PhD thesis of 1947). In that work they used perylene and a few other fluorescent dyes. A few years later Meir Shinitzky (who had been a postdoctoral fellow with Weber and a coauthor on the 1971 paper) and Yechezkel Barenholz introduced the probe diphenylhexatriene (DPH), which arguably became the most popular fluorescence polarization membrane probe of all time.

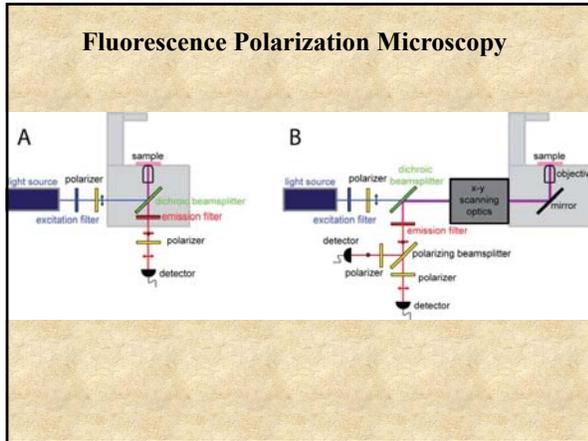


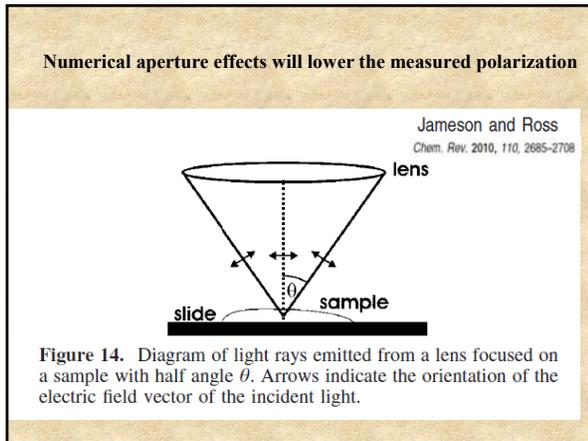


Anisotropy of DPH in DPPC vesicles as a function of temperature (courtesy of Ivo Konopasek from <http://web.natur.cuni.cz/~konop/gallery.php>)

**Fluorescence polarization
platereaders are available for
High Throughput Screening
experiments**







CARBOFLUORESCIN DYE ORIENTATION IN RED CELL MEMBRANE STUDIED BY MICROSCOPIC FLUORESCENCE POLARIZATION

DANIEL AVALOS, Biophysics Research Division and Department of Physics, The University of Michigan, 460-460B Michigian 48109 U.S.A.

Biophys.J. 1979 Jun;26(3):557-73

APPENDIX

High Aperture Correction of Fluorescence Polarization

We must first define two right-handed coordinate systems (Fig. 7): one (the X system) is a microscope-fixed reference frame; the other (the X' system) is a nonfixed frame convenient for analyzing the polarization of rays in object space. Both of these systems share the same origin at which is located the transition dipole moments of the fluorescent sample under view in the object space of the objective. We assume for simplicity of units that the absorption and emission dipole moments are of unit magnitude.

In the X system, X_1 is parallel to the optical axis of the microscope and the X_2 axis is parallel to the electric field vector of the polarized excitation light. An emission dipole has components along these axes denoted by (x_1, x_2, x_3) . A typical emitted fluorescence ray propagates toward the objective in a direction given by polar angle, α , and azimuthal angle, ϕ , with respect to the X_1 and X_2 axes, respectively. Because

FIGURE 7. Coordinate systems X (microscope-fixed) and X' (emitted ray-fixed). The fluorescence sample is at origin O . The meridional plane is shaded. The emitted ray travels along X'_1 in object space and then parallel to X_1 in image space. Axis X_1 is the direction of the excitation light polarization. Axes X_2 and X'_2 both make the same angle θ with the meridional plane.

